

# Frequency-domain self-force calculations using Gegenbauer reconstruction

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# Gravitational scattering

## Why study scattering?

- Clean, well-defined asymptotic in/out states.
- Probe strong-field (sub-ISCO) at low energy.
- EOB Hamiltonian  $H_{\text{EOB}}$  completely determined by scatter angle  $\chi(E, L)$ . [Damour 2016+]
- B2B maps between scatter and bound orbit. [Kalin & Porto 2020+]

## Complementary approaches include:

- **Numerical relativity:** low-separation, short durations,  $\approx$  equal mass.
- **Post-Minkowskian expansion (expansion in  $G$ ):** weak-field, arbitrary mass-ratio, analytical.
- **Gravitational self-force (expansion in mass-ratio  $\epsilon$ ):** strong- and weak-field, small mass-ratio, numerical.

Interest in developing numerical GSF calculations along scatter orbits.

# Scalar-field self-force and scattering

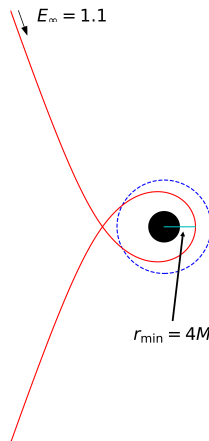
- Scalar charge  $q$  with mass  $\mu$  scattered off a Schwarzschild black hole of mass  $M$ . Scalar field:

$$\nabla\Phi = -4\pi q \int \frac{\delta^4(x^\alpha - x_p^\alpha(\tau))}{\sqrt{-g}} d\tau,$$

where  $\epsilon := q^2/\mu M \ll 1$  is the expansion parameter.

- At leading order take  $x_p^\alpha(\tau)$  to be a scatter geodesic: parameterised by energy at infinity,  $E_\infty$ , and periapsis radius  $r_{\min}$ .
- Particle feels a self-force due to interaction with its own scalar field:

$$u^\beta \nabla_\beta (\mu u^\alpha) = q \nabla_\beta \Phi^R := F^\alpha.$$



# Numerical self-force calculations

Decompose the field into spherical harmonics centred on the Schwarzschild black hole:

$$\Phi(t, r, \theta, \phi) = \frac{1}{r} \sum_{\ell=0}^{\infty} \sum_{m=-\ell}^{+\ell} \psi_{\ell m}(t, r) Y_{\ell m}(\theta, \phi)$$

Self-force can be computed using mode-sum regularisation:

$$F_{\alpha}(\tau) = \sum_{\ell=0}^{\infty} \left\{ \nabla_{\alpha} \left[ \frac{1}{r} \sum_{m=-\ell}^{+\ell} \psi_{\ell m} Y_{\ell m} \right]_{x_p^{\pm}(\tau)} \mp \left( \ell + \frac{1}{2} \right) \underbrace{A_{\alpha}(\tau) - B_{\alpha}(\tau)}_{\text{regularisation parameters}} \right\}$$

## Time-domain

- Solve  $(1+1)d$  **PDEs** for  $\psi_{\ell m}(t, r)$ .
- Scattering: [Long & Barack 2209.03740]

## Frequency-domain

- Construct  $\psi_{\ell m}(t, r)$  from frequency-modes, obtained by solving **ODEs**.
- Scattering: [Whittall & Barack 2305.09724]

# Direct reconstruction

- Frequency-modes  $\psi_{\ell m \omega}$  obey ODE

$$\frac{d^2 \psi_{\ell m \omega}}{dr_*^2} - [V_\ell(r) - \omega^2] \psi_{\ell m \omega} = S_{\ell m \omega}(r)$$

- Retarded solution expressed in terms of homogeneous solution basis  $\psi_{\ell \omega}^\pm(r)$  using variation of parameters:

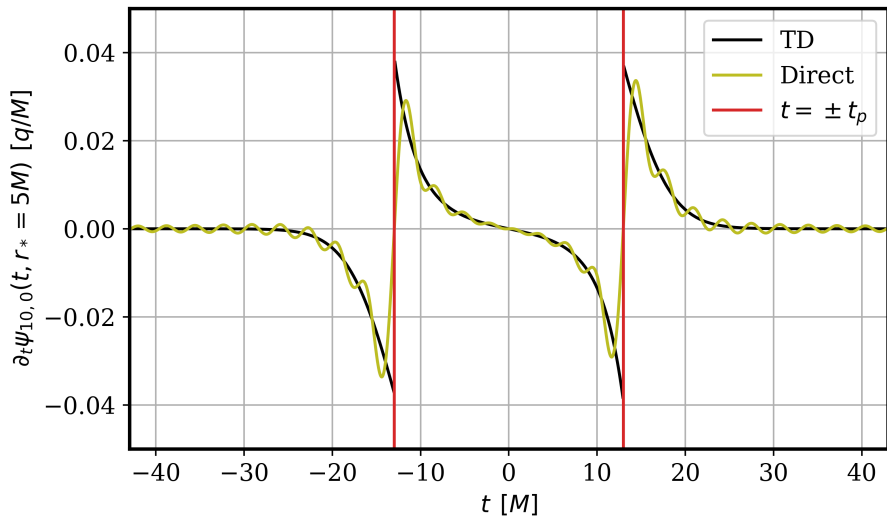
$$\psi_{\ell m \omega}(r) = \psi_{\ell \omega}^+(r) \int_{r_{\min}}^r \frac{\psi_{\ell \omega}^-(r') S_{\ell m \omega}(r')}{W_{\ell \omega} f(r')} dr' + \psi_{\ell \omega}^-(r) \int_r^{+\infty} \frac{\psi_{\ell \omega}^+(r') S_{\ell m \omega}(r')}{W_{\ell \omega} f(r')} dr'$$

- Direct reconstruction:

$$\psi_{\ell m}(t, r) \approx \Psi_{\ell m}(t, r; \omega_{\max}) := \int_{-\omega_{\max}}^{+\omega_{\max}} \psi_{\ell m \omega}(r) e^{-i\omega t} d\omega$$

is **not practical** due to the **Gibbs phenomenon**.

# Direct reconstruction



# EHS reconstruction [Barack, Ori & Sago 2008]

- **EHS reconstruction**: recover  $\psi_{\ell m}(t, r)$  separately in  $r \leq r_p(t)$  and  $r \geq r_p(t)$  using homogeneous solutions.
- For example, field modes in the “internal” region  $r \leq r_p(t)$  reconstructed from

$$\tilde{\psi}_{\ell m \omega}^{-}(r) := \psi_{\ell \omega}^{-}(r) \int_{r_{\min}}^{+\infty} \frac{\psi_{\ell \omega}^{+}(r') S_{\ell m \omega}(r')}{W_{\ell \omega} f(r')} dr'.$$

- Restores **exponential, uniform convergence**.
- **External reconstruction**: EHS cannot be applied in  $r > r_p(t)$  for scatter orbits.
- **Cancellation problem**: significant loss of precision at high eccentricities [van de Meent 2016]. Limited to  $r_p \sim r_{\min}$  for scatter orbits.

# Gibbs-complementary reprojection

- Say that a basis of functions  $\{C_k\}$  (on a compact time interval  $I$ ) form a **Gibbs-complement** to the Fourier basis if: [Gottlieb & Shu 1998]
  - ① for any function  $f$  which is analytic on  $I$ , the expansion in the  $\{C_k\}$  basis converges exponentially to  $f$ , **and**
  - ② the projection of the high-frequency Fourier content onto the low-degree  $C_k$  can be made exponentially small.
- **Gegenbauer polynomials**  $\{C_k^\lambda(s)\}$  orthogonal on  $[-1, 1]$  wrt weighted inner product,

$$\int_{-1}^{+1} (1 - s^2)^{\lambda-1/2} C_n^\lambda(s) C_m^\lambda(s) ds = h_n^\lambda \delta_{nm}.$$

- ▶ Generalisation of Legendre ( $\lambda = 1/2$ ) and Chebyshev ( $\lambda = 0, 1$ ) polynomials.
- ▶ **Gibbs-complementary** to the Fourier basis.



# Gegenbauer reconstruction [Gottlieb & Shu 1992 (et al), 1994, 1995, 1997]

Suppose  $\psi_{\ell m}(t, r)$  is analytic (at fixed  $r$ ) on interval  $a \leq t \leq b$ :

- 1 Compute the partial Fourier integrals using the inhomogeneous modes  $\psi_{\ell m \omega}(r)$  for  $t \in [a, b]$ :

$$\Psi_{\ell m}(t, r; \omega_{\max}) := \int_{-\omega_{\max}}^{+\omega_{\max}} \psi_{\ell m \omega}(r) e^{-i\omega t} d\omega.$$

- 2 Project onto the Gegenbauer basis:

$$g_k^\lambda(r; \omega_{\max}) := \frac{1}{h_k^\lambda} \int_{-1}^1 (1-s^2)^{\lambda-1/2} \Psi_{\ell m}(t(s), r; \omega_{\max}) C_k^\lambda(s) ds.$$

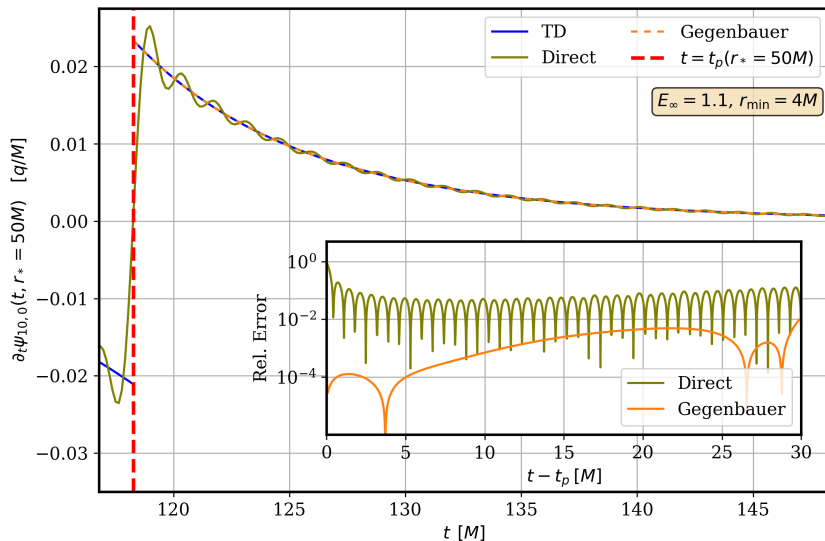
where  $t(s) = [(b-a)s + (a+b)]/2$ .

- 3 Approximate

$$\psi_{\ell m}(t, r) \approx \sum_{k=0}^N g_k^\lambda C_k^\lambda(s(t)).$$

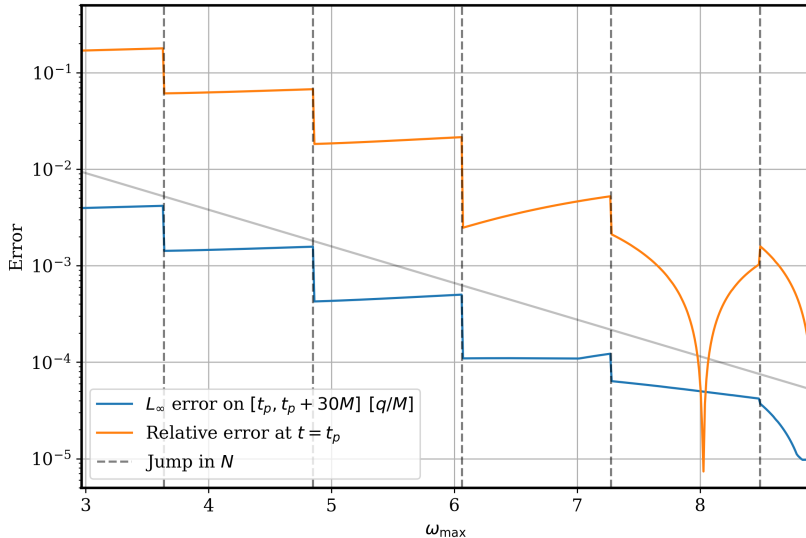
Gegenbauer approximant converges **uniformly** and **exponentially** on  $a \leq t \leq b$ , provided  $N, \lambda$  and  $\omega_{\max} \rightarrow \infty$  in **linear proportion**.

# Internal reconstruction: $r \leq r_p(t)$



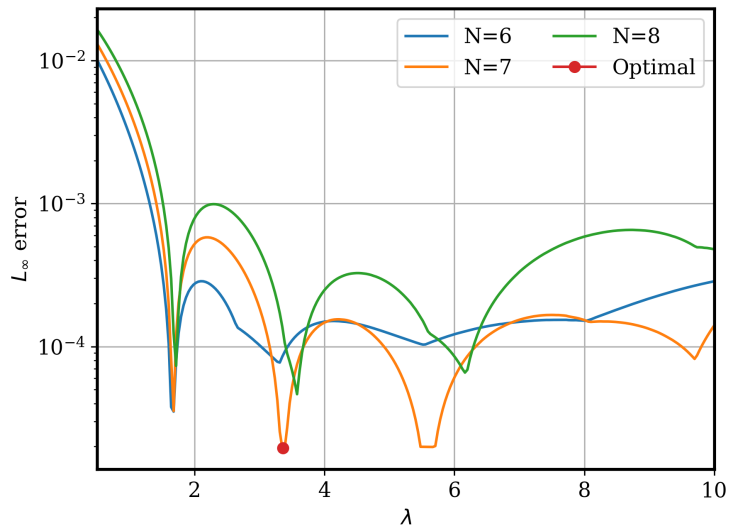
Gegenbauer reconstruction effective and outperforms Direct reconstruction.

# Convergence



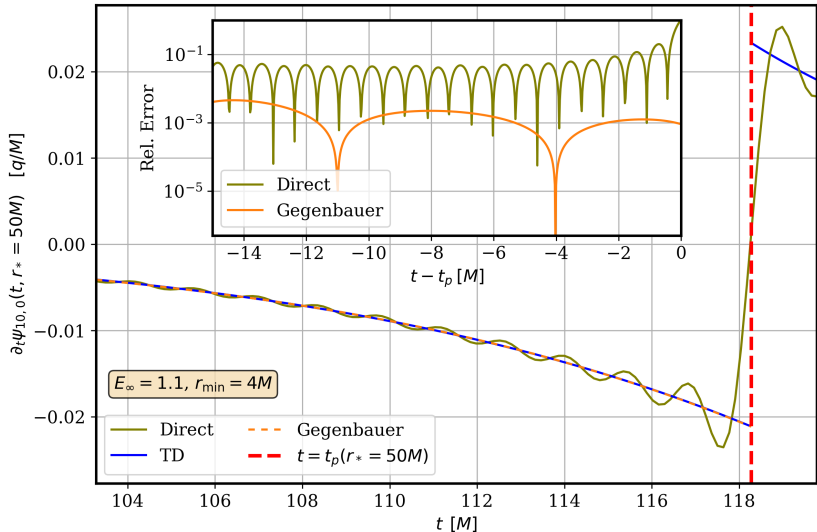
Convergence expected to be uniform and exponential.

# Choice of $N$ and $\lambda$



Parameters  $N$  and  $\lambda$  must be chosen appropriately.

## External reconstruction: $r \geq r_p(t)$



Gegenbauer reconstruction enables calculations in  $r > r_p(t)$ .

# Concluding remarks

- **Coming soon to arXiv:** proof of concept calculations for scalar-field modes  $\psi_{\ell m}(t, r)$  with scattering source.
- Working towards full scalar-field self-force scatter calculation:
  - ▶ **Computational efficiency:** reconstruction, radial integrals in FD.
    - TD reconstruction cheaper than FD mode calculation.
    - Economise FD calculation by reusing integrals/integrands, or use ODE integration.
  - ▶ **Optimal choices of  $N/\omega_{\max}$  and  $\lambda/\omega_{\max}$ .**
    - Calibrate to EHS then update using convergence tests?
    - Choices unlikely to be optimal, but may be adequate.
- **Gravity** is the ultimate aim .
- Advanced reprojection methods?
  - ▶ Alternative bases (e.g. Freud polynomials [Gelb & Tanner 2006]).
  - ▶ Inverse polynomial reconstruction [Shizgal & Jung 2003]
  - ▶ ... and more! e.g. [Adcock & Hansen 2011]
- New applications?
  - ▶ Easy to implement as post-processing step.